

Non-local NC simulations in stellarators/heliotrons using the FORTEC-3D code

Presented by J.L. Velasco¹,
including work and discussion by S. Satake², A. Dinklage³, M. Yokoyama², D. López-Bruna¹ and many people from TJ-II team¹, LHD Exp. Group² and W7-AS Team³ and others^{4,5,6}

¹ *Laboratorio Nacional de Fusión, CIEMAT, Madrid, Spain.*

² *National Institute for Fusion Science, Toki, Japan.*

³ *Max-Planck Institute fuer Plasmaphysik, Greifswald, Germany.*

⁴ *Research Organization for Information Science and Technology, Japan.*

⁵ *Department of Nuclear Engineering, Kyoto University, Kyoto, Japan.*

⁶ *Princeton Plasma Physics Laboratory, Princeton, NJ, USA*

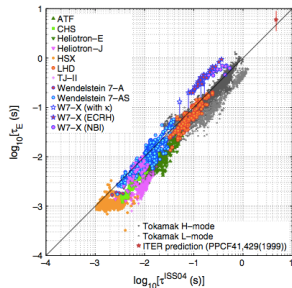
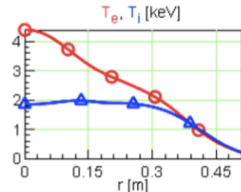
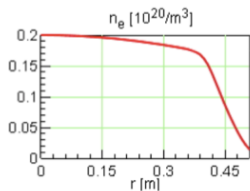
Gyrokinetic Theory Working Group Meeting, Madrid, 2014

Questions that we want to answer

Long term goal: support predictive simulations of W7-X scenarios and stellarator reactors (based on NC transport in the core + simplified turbulent transport in the edge + sources).

$$\frac{3}{2} \frac{\partial n_b T_b}{\partial t} + \frac{1}{V'} \frac{\partial}{\partial r} V' Q_b = P_b.$$

[Pedersen EPS 2014, Dinklage NF 2007]



- How accurately does NC theory describe radial heat and particle transport in a stellarator? / Where do we need GK simulations?
- Does including higher order NC terms improve the description?
- Is the situation different for particle and heat transport?

What to expect

- The long-wave-length radial electric field is expected to be set by Neoclassics in stellarators far enough from quasisymmetry e.g. [Calvo PPCF 2013] because:

$$\langle \mathbf{J} \cdot \nabla r \rangle = \langle \mathbf{J}_{NC} \cdot \nabla r \rangle + O(\rho^*{}^3), \quad \rho^* \equiv \frac{\rho_L}{L_0},$$

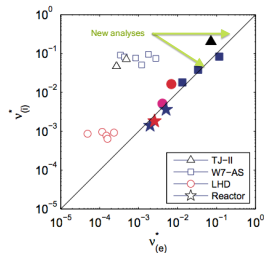
but no such result exists for the radial particle and energy fluxes.

- Empirical (and tautological) rule: neoclassical transport dominates when it is large enough:
 - Low collisionalities (core region, where T_i and T_e are large).
 - Configuration far from omnigenity ($1/\nu$) [Parra submitted 2014].
- ρ^* is small (e.g., $\rho^* \sim 10^{-2}$ for TJ-II) so higher-order NC terms should be negligible.
 - High T_i and small $|E_r|$ is a standard situation in the core region of stellarators. $\Rightarrow \mathbf{v}_M \cdot \nabla \delta f$ may be non-negligible for ion transport.

Previous results

[Dinklage NF 2013]

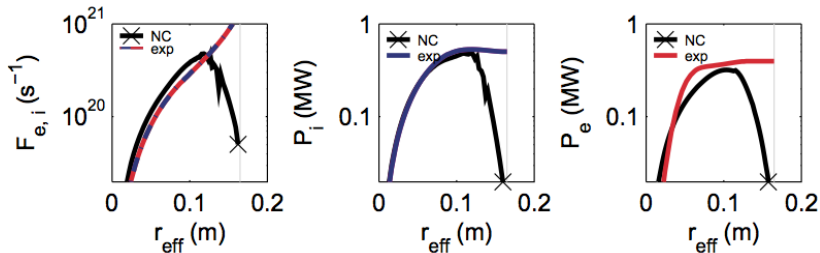
- Medium-size stellarator/heliotrons at medium-to-high densities, $n_e > 4 \times 10^{19} m^{-3}$, $T_b \sim keV$.
- Ion root (E_r small and negative), ion energy transport is dominant.



- Predicted NC energy fluxes roughly comply with experimental fluxes (from transport balance) in the core region (but there is room for improvement).
- Bad agreement in outer positions.
- Preliminary results of simulations with FORTEC-3D predict that non-local effects have more impact in Q_b than in Γ .

[Satake EPS 2014]

[Dinklage NF 2013] W7-AS #34609



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[Satake EPS 2014]

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Solve the time-independent drift kinetic equation for $\delta f(r, \zeta, \theta, v, \xi)$:

$$\left[(\mathbf{v}_{\parallel} + \mathbf{v}_{\mathbf{E}} + \mathbf{v}_{\mathbf{M}}) \cdot \nabla + \dot{v} \frac{\partial}{\partial v} + \dot{\xi} \frac{\partial}{\partial \xi} \right] \delta f - C(\delta f) = -(\mathbf{v}_{\mathbf{M}} \cdot \nabla + \frac{\partial}{\partial K}) f_M + \mathcal{P} f_M.$$

Differences w.r.t. local and monoenergetic approximation (DKES, [Beidler NF 2011]):

- $\mathbf{v}_{\mathbf{E}} = \frac{\mathbf{E} \times \mathbf{B}}{B^2}$ is *compressible*.
- $C(\delta f) + \mathcal{P} f_M$ includes pitch-angle and energy scattering + momentum conservation.
- $\mathbf{v}_{\mathbf{M}} = \frac{v^2(1+\xi^2)}{q} \frac{\mathbf{B} \times \nabla B}{B^3}$ is the magnetic drift \Rightarrow need to keep $\frac{\partial \delta f}{\partial v}$.

Take moments:

$$\Gamma_b = \left\langle \int d^3v \delta f v_M \cdot \nabla r \right\rangle, \quad Q_b = \left\langle \int d^3v v^2 v_M \delta f \cdot \nabla r \right\rangle.$$

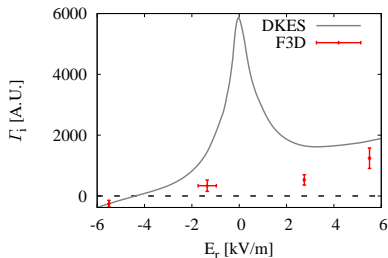
E_r from ambipolarity of NC fluxes: $\Gamma_e(E_r) = \Gamma_i(E_r)$.

The $\mathbf{v}_M \cdot \nabla \delta f$ term

Solve the time-independent drift kinetic equation for $\delta f(r, \zeta, \theta, v, \xi)$:

$$\left[(\mathbf{v}_{\parallel} + \mathbf{v}_E + \mathbf{v}_M) \cdot \nabla + v \frac{\partial}{\partial v} + \xi \frac{\partial}{\partial \xi} \right] \delta f - C(\delta f) = -(\mathbf{v}_M \cdot \nabla + \frac{\partial}{\partial K}) f_M + \mathcal{P} f_M.$$

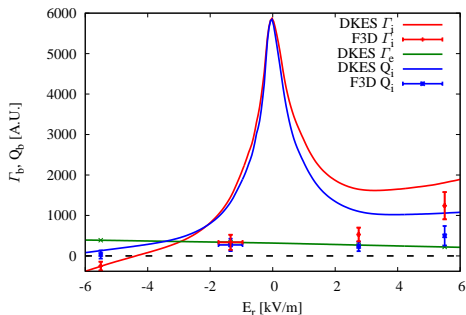
- $\mathbf{v}_E \cdot \nabla \delta f$ is formally higher order, but need to keep it for low collisionalities.
- $\mathbf{v}_M \cdot \nabla \delta f$ of same order than $\mathbf{v}_E \cdot \nabla \delta f$.
 - Usually neglected.
 - For $|E_r|$ small enough may be a bad idea.



Overestimate of NC radial fluxes in the core of stellarators? \Rightarrow
 \Rightarrow Underestimate of turbulent fluxes?

Effect on theoretical predictions

- $\mathbf{v}_M \cdot \nabla \delta f$ reduces Γ_i and Q_i for a given E_r .
- E_r from ambipolarity of NC fluxes: $\Gamma_e(E_r) = \Gamma_i(E_r)$
- Γ_e depends (more) weakly on E_r .



- E_r slightly less negative.
- Γ very slightly reduced.
- Larger reduction in Q_i .

Parameter dependence

The size of the peak is larger:

- The smaller the collisionality.
- The larger the deviation from omnigenity.

(competing effects).

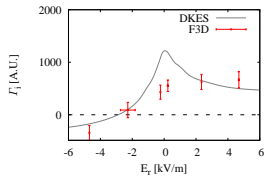
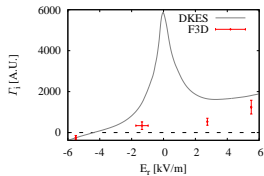
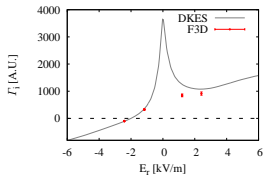
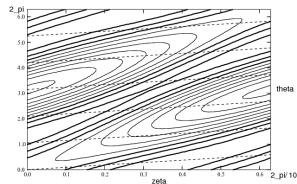
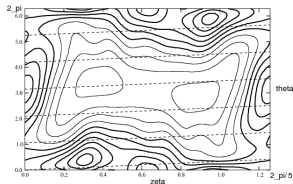
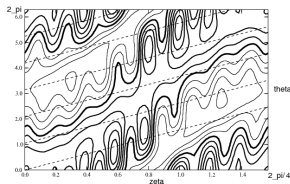
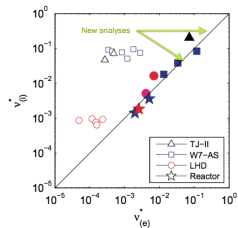
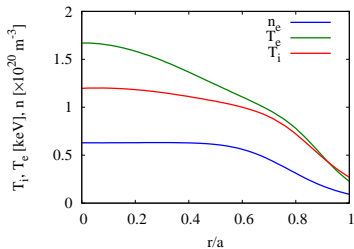
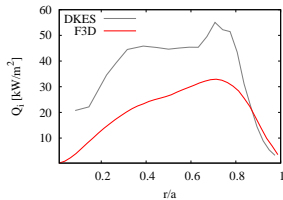
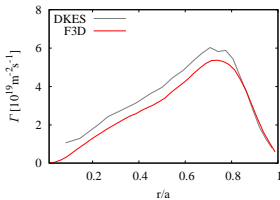
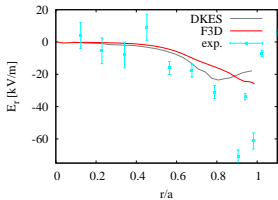


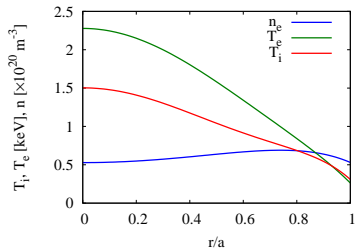
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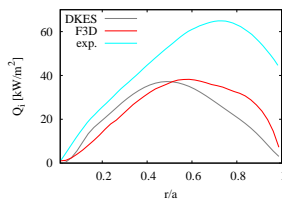
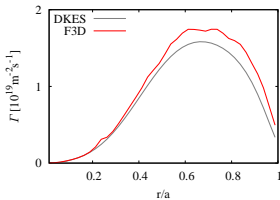
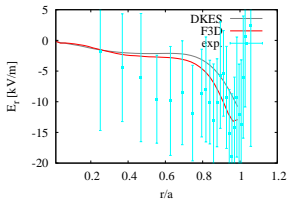


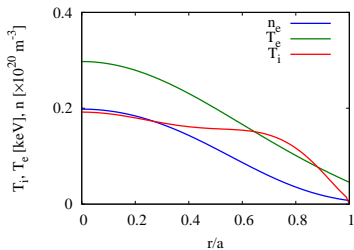
- Low-collisionality W7-AS case.
- Reasonable agreement in E_r .
- Q_i reduced by a factor 2.





- Lowest collisionality and ripple.
- Q_i from 25% below exp. to 50% below.





- Agreement in E_r and Γ .
- Predicted Q_i always below experiment (high collisionality).

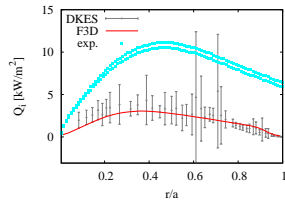
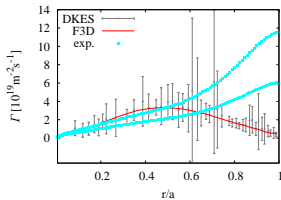
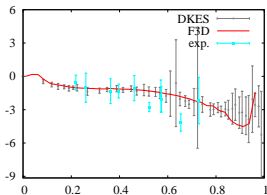


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- Neglecting higher order NC terms can lead to underestimate of radial energy flux in the core of stellarators, specially:
 - if the stellarator is not optimized.
 - if the collisionality is small.
- Need to calculate turbulent contribution.
- Bootstrap current? (time-consuming!).
- When W7-X starts operation, check NC predictions.