

# Some features of RF discharge development in the Uragan-3M torsatron

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## INTRODUCTION

In the  $\ell=3/m=9$  U-3M torsatron with an open helical divertor a hydrogen plasma is RF produced and heated ( $\omega \lesssim \omega_{ci}$ ). The RF power is introduced into the plasma with using an unshielded frame-like antenna with a broad spectrum of parallel wavelengths.

**O.M. Shvets., I.A. Dikij, S.S. Kalinichenko et al. Nucl. Fusion 1986 26 23.**

**M.D. Carter, A.I. Lysoivan, V.E. Moiseenko et al. Nucl. Fusion 1990 30 723**

**Yu.G. Zalessky, P.I. Kurilko, N.I. Nazarov et al. Fizika plazmy 1989 15 1424 (in Russian).**

Continuation of works:

**V. V. Chechkin, L. I. Grigor'eva, Ye. L. Sorokovoy et al. 2009 Plasma Phys. Repts. 35 852.**

**I. M. Pankratov, A. A. Beletskii, V. L. Berezhnij et al. 2010 Contrib. Plasma Phys. 50 520.**

Time evolution is studied of:

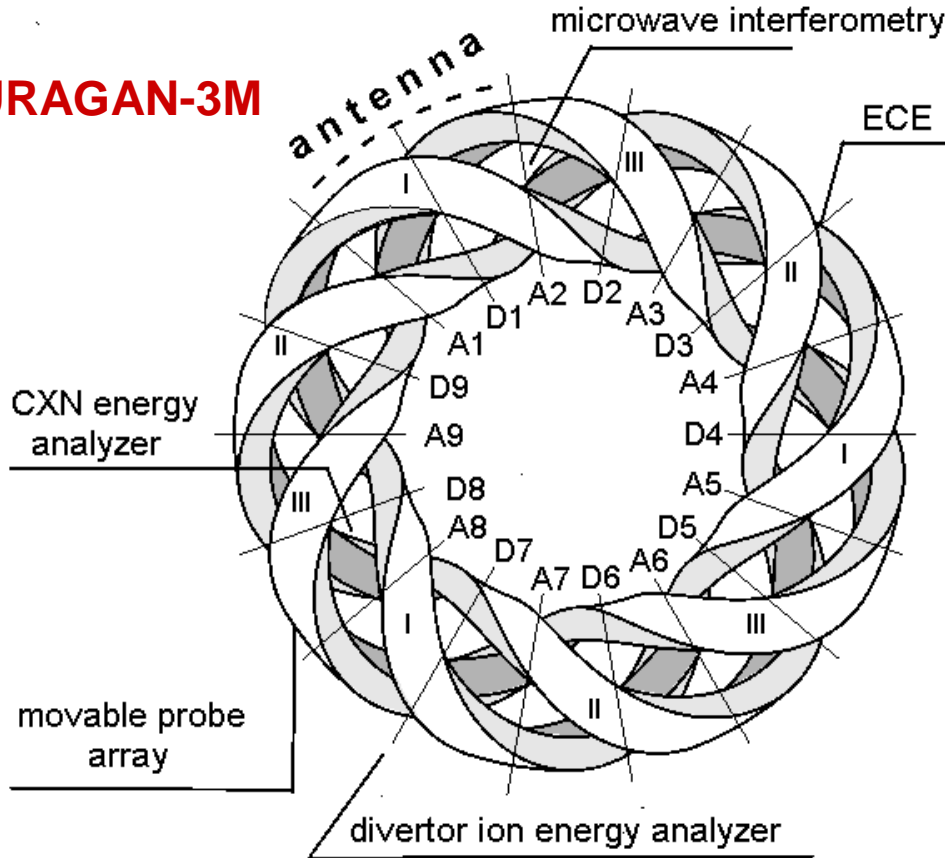
- Density  $\bar{n}_e$  and electron cyclotron emission (ECE) from the central plasma at different values of RF power  $P$  fed to the antenna;
- Edge radial electric field  $E_r$  and edge turbulent transport;
- Fast ion (FI,  $\gtrsim 500$  eV) generation and loss.

The results are of interest to:

- Target plasma formation for subsequent production and heating of a denser plasma using another, shorter-wave antenna;
- Studying processes which are determinative in formation of the H-like mode states in U-3M and need a more detailed investigation.

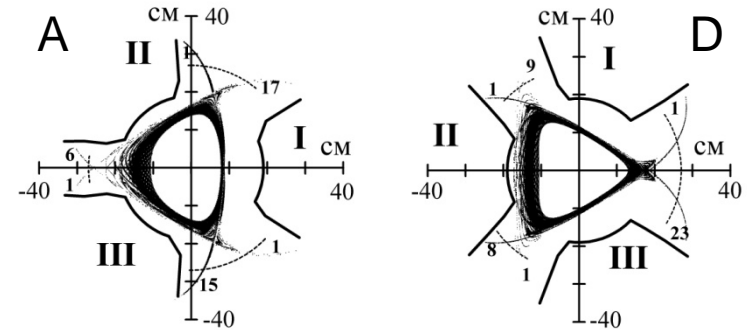
# EXPERIMENTAL CONDITIONS

**URAGAN-3M**



I, II, III, helical coils

A1, D1, A2, D2, ..., A9, D9,  
symmetric poloidal cross-  
sections in the periods 1, 2, ..., 9



RF antenna is mounted under  
coils I and III with the leads in D1

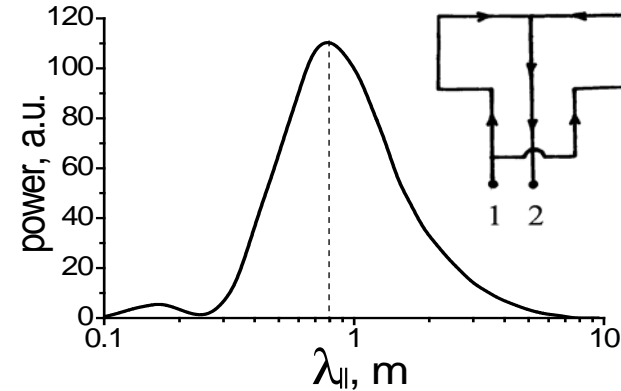
$\ell = 3 / m = 9$  torsatron,  $R_0 = 100$  cm,  $\bar{a} \approx 12$  cm,  $\iota(a)/2\pi \approx 0.3$ .

Helical coils only,  $B_\phi \lesssim 1$  T, with the whole magnetic system enclosed into a large ( $70 \text{ m}^3$ ) vacuum chamber – an open helical divertor.

Continuous fuelling gas (hydrogen) admission,  $p \sim 10^{-5}$  Torr.

RF plasma production and heating,  $\omega \lesssim \omega_{ci}$  under local Alfvén resonance (LAR) conditions:  $N_{\parallel}^2 = \varepsilon_1$  ( $N_{\parallel} = k_{\parallel} c / \omega$ ,  $\varepsilon_1 \approx \omega_{pi}^2 / (\omega_{ci}^2 - \omega^2)$ ).

### Twisted frame-like antenna.



$\lambda_{\parallel} \sim 40-400$  cm;  $n_i \approx (7-0.07) \times 10^{12}$  cm $^{-3}$ ;  
 $\lambda_{\parallel \max P} \approx 80$  cm;  $n_{i \max P} \approx 1.8 \times 10^{12}$  cm $^{-3}$ .

**Yu.G. Zalessky, P.I. Kurilko, N.I. Nazarov et al.**  
*Fizika plazmy* 1989 15 1424 (in Russian).

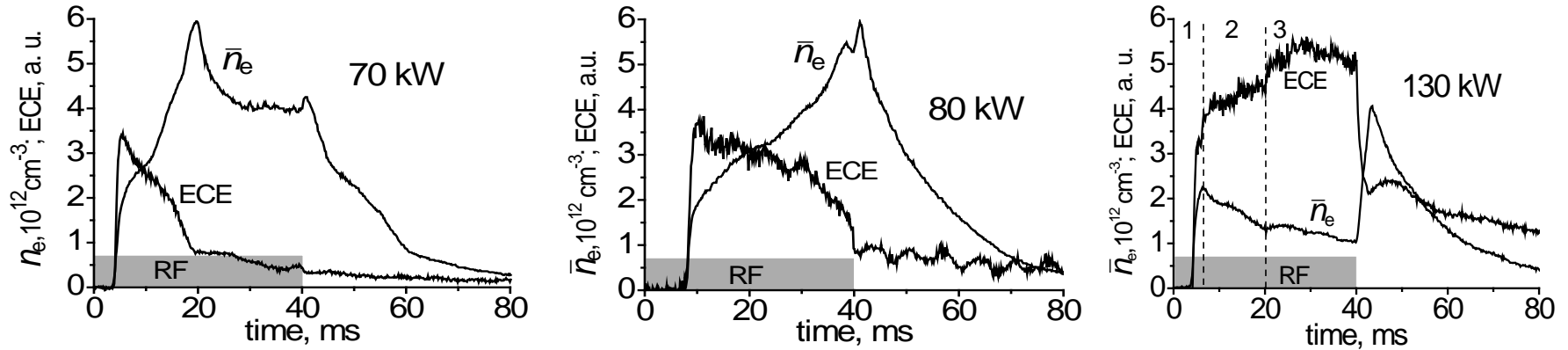
$B_{\phi} = 0.72$  T;  $\omega / 2\pi = 8.8$  MHz ( $\omega = 0.8\omega_{ci}(0)$ )  
 RF power fed to the antenna  $P \lesssim 150$  kW

$\bar{n}_e \sim \text{units} \times 10^{12}$  cm $^{-3}$

Two groups of ions:  $T_{i1} \sim$  tens eV (optical spectroscopy);  $T_{i2} \sim 300-600$  eV (CXN)

$T_e(0) \sim 500-700$  eV (ECE).

# TIME EVOLUTION OF ELECTRON DENSITY AND RADIATION TEMPERATURE AT DIFFERENT LEVELS OF RF POWER



**$P \sim 60 \rightarrow 80$  kW:**  $\bar{n}_e$  attains  $\sim 2 \times 10^{12} \text{ cm}^{-3}$  (in  $\sim 2.5$  ms at  $P=60$  kW), then the density rise is slowed down ( $\bar{n}_e(t)$  bend), with maximum  $\bar{n}_e \approx 6 \times 10^{12} \text{ cm}^{-3}$  attained the later the higher  $P$  is. The central ECE attains its maximum at  $\bar{n}_e \approx 2 \times 10^{12} \text{ cm}^{-3}$ , then multiply drops, presumably indicating a LAR zone shift to the periphery.

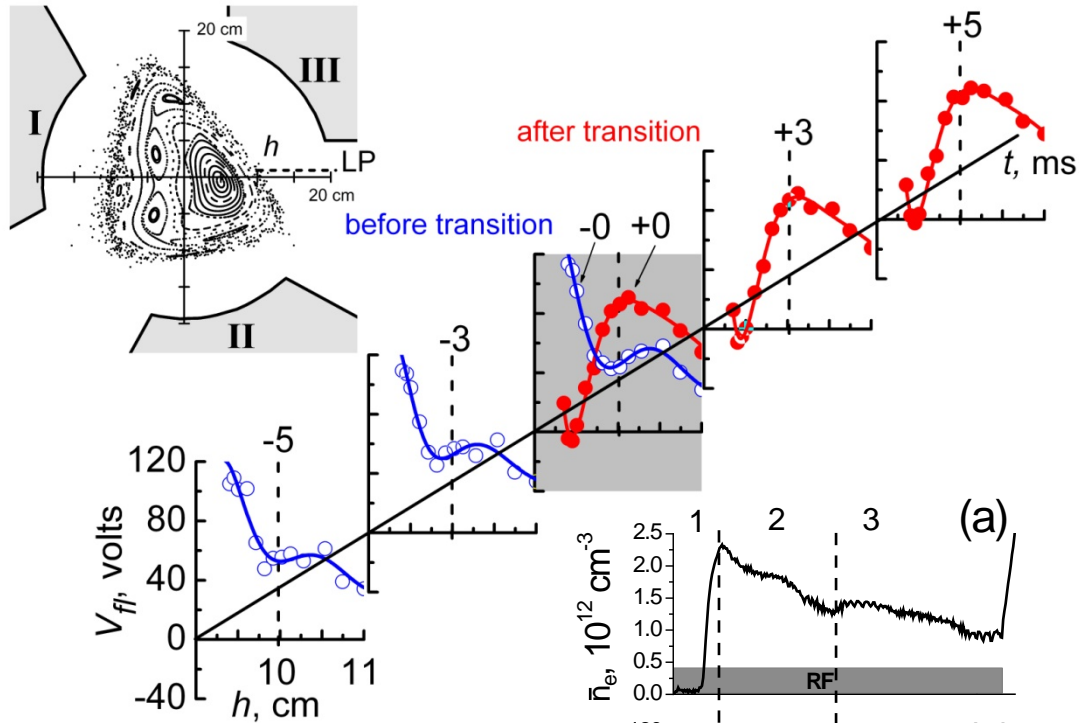
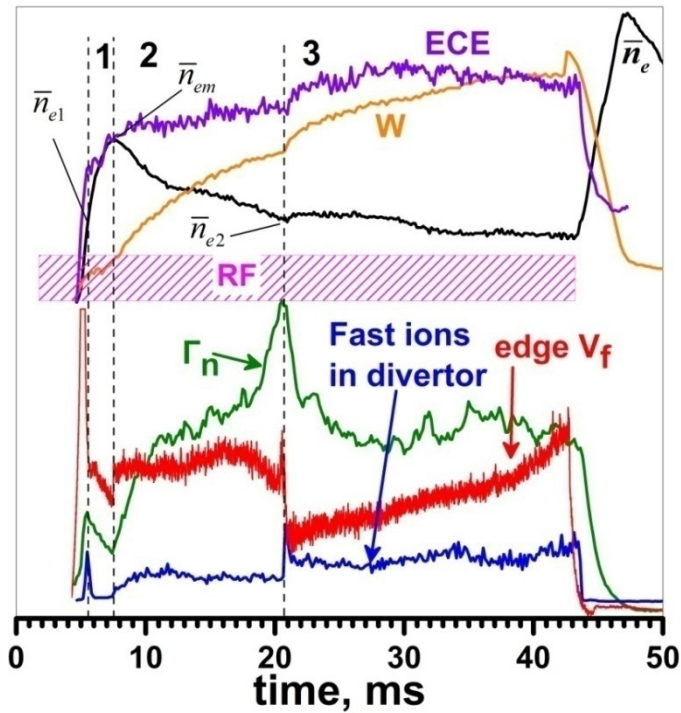
**$P > 80$  kW:**  $\bar{n}_e$  does not attain  $\approx 6 \times 10^{12} \text{ cm}^{-3}$ , starting to decay before RF off, and the faster the higher  $P$  is, this being a manifestation of confinement degradation with power. At  $P \gtrsim 100$  kW the maximum density becomes  $\bar{n}_{em} \approx 2 \times 10^{12} \text{ cm}^{-3}$  ( $\bar{n}_e(t)$  bend), while the high level of ECE persists over all the RF pulse, indicating optimum RF heating conditions.

**$P > 110$  kW,** as soon as the slowly decaying density approaches  $\bar{n}_{e2} \approx 1.2 \times 10^{12} \text{ cm}^{-3}$ , some commonly observed indications of the H-like mode transition occur:

The  $\bar{n}_e$  rise after RF off is due to the transport reduction with plasma cooling at a continuous influx of the fuelling gas (confirmed by numerical modeling).

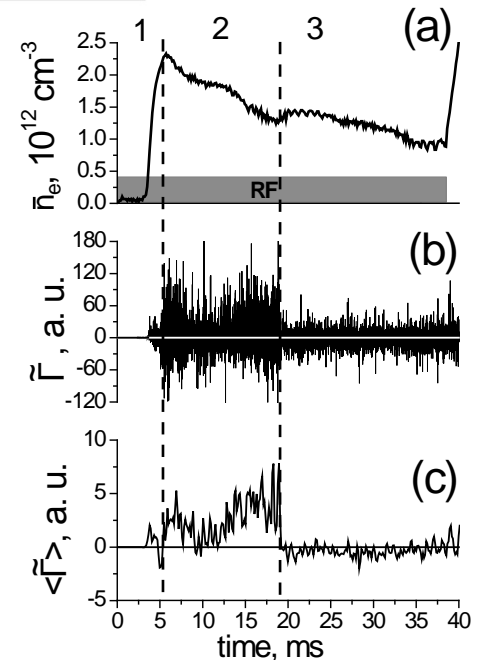
**N.T. Besedin et al. 1993 Stellarators and Other Helical Confinement Systems. Collection of Papers Presented at the IAEA TCM (Garching, Germany, 1993) (Vienna: IAEA) p. 277**

$P = 130 \text{ kW}$



### Indications of H-like mode transition:

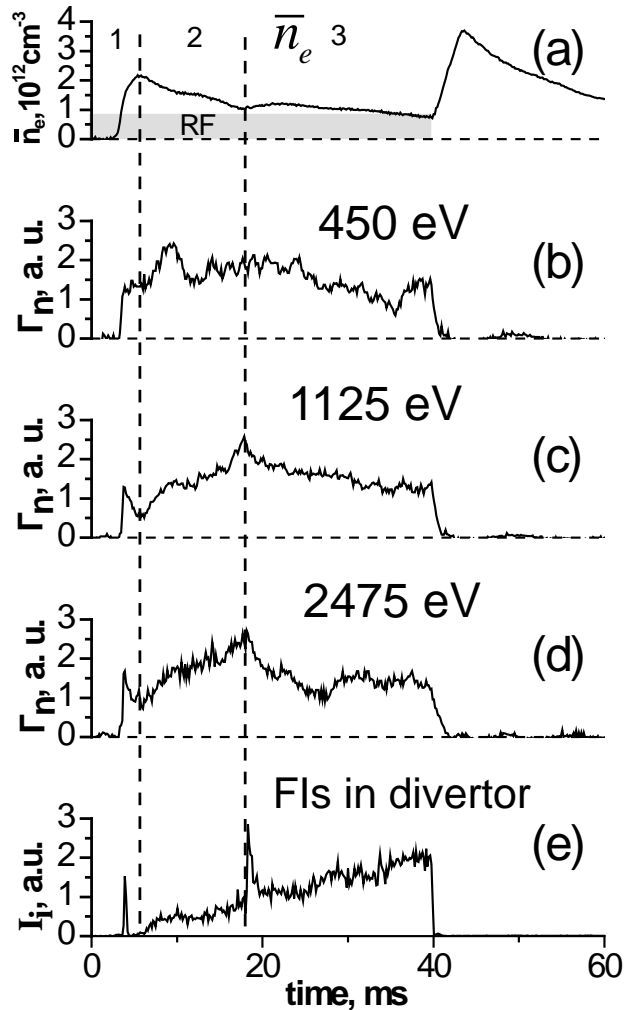
- Temporal suspension of  $\bar{n}_e$  decay and even some  $\bar{n}_e$  rise;
- $W_{\text{dia}}(t)$  and  $ECE(t)$  bendings toward speeding-up of their rise;
- Step-like change in the edge potential  $V_f$  caused by an edge  $|E_r|$  (and  $E_r$  shear) rise;
- Edge turbulence suppression



A distinct time correlation between the level of edge turbulent flux  $\tilde{\Gamma}$  and the rate of average density decay

# TIME BEHAVIOR OF FAST IONS ( $T_{i2}$ GROUP, $\approx 500$ eV) AND THEIR LOSS.

## The link between FI loss and H-mode transition



(b):  $W_{\perp} < 500$  eV: no correlation with  $\bar{n}_e$  (t)

(c, d):  $W_{\perp} > 500$  eV:  $\Gamma_n$  (t) changes in antiphase with  $\bar{n}_e$ ; resonance-like dependence of  $\Gamma_n$  on  $\bar{n}_e$

(e): a short-time ( $< 0.5$  ms) single burst of FI outflow to the divertor ("burst of FI loss") just before H-like transition – an indication of transition to be triggered by the FI loss.



## SUMMARY

1. Using the frame-like antenna with a broad spectrum of parallel wavelengths  $\lambda_{\parallel}$ , a cold plasma with  $\bar{n}_e \sim (4-7) \times 10^{12} \text{ cm}^{-3}$  ( $P \sim 60-80 \text{ kW}$ ) and a plasma with high  $T_e(0)$  and  $\bar{n}_e \sim 2 \times 10^{12} \text{ cm}^{-3}$  ( $P > 100 \text{ kW}$ ) can be produced in U-3M. At present both regimes are being investigated as possible sources of a target plasma for production and heating of denser (up to  $\sim 10^{13} \text{ cm}^{-3}$ ) plasmas using another, shorter-wave antenna with azimuthal currents.

Kasilov S.V., Lysojvan A.I., Moiseenko V.E., Plyusnin V.V. *Stellarators and Other Helical Confinement Systems* (Vienna: IAEA, 1993, p. 447).

2. Processes occurring during the RF pulse can be divided into slower and faster ones by their characteristic time scale.

**Slower processes** (tens/units ms):  $\bar{n}_e$  decay and energy content  $W_{\text{dia}}$  rise (both  $\sim 10$  ms); FI concentration rise (units ms). As a result, a peculiar “coupling resonance” sets in where optimum conditions for FI generation at  $\bar{n}_e \approx 1.2 \times 10^{12} \text{ cm}^{-3}$  are realized.

**Faster processes** (hundreds/tens  $\mu\text{s}$ ) responsible for the H-like mode transition in itself: the burst of FI loss (hundreds  $\mu\text{s}$ )  $\rightarrow$  the edge  $E_r$  bifurcation toward a more negative value with the higher  $E_r$  shear (tens  $\mu\text{s}$ )  $\rightarrow$  suppression of the edge turbulence and the turbulence-induced transport (tens  $\mu\text{s}$ ) with improvement of plasma confinement as a final result.

**$\tau_E$  is shown to increase from 2.7 ms to 5.1 ms**

V.K. Pashnev, et al. *Problems At. Sci. Tech.* 2010 No. 6 Series: Plasma Phys. (16) p. 17



3. Similar to other devices, in U-3M a specific radial profile of the edge electrostatic potential with the “well” near the LCFS and a lower  $E_r$  shear are formed in the L-like state and can be attributed, by analogy, to the non-ambipolar ion orbit loss. The previously formed  $E_r$  shear considerably increases with transition, with the potential “well” shifting inside the LCFS.

4. Triggered by a short-time enhanced FI loss, the state of discharge with the modified edge  $E_r$  profile where the turbulent transport is suppressed persists for a comparatively long period without recovering the pre-transitional high level of the edge turbulence.

In more detail see

*PAST (Problems of Atomic Science and Technology). Series “Plasma Physics”. 2012. No. 6 (82), p. 3*

V. V. Chechkin, I. M. Pankratov, L. I. Grigor’eva, A. A. Beletskii, A. A. Kasilov, Ye. D. Volkov, V. Ye. Moiseenko, V. S. Voitsenya, V. K. Pashnev, P. Ya. Burchenko, A. V. Lozin, S. A. Tsybenko, A. S. Slavnyj, M. Dreval, A. P. Litvinov, A. Ye. Kulaga, R. O. Pavlichenko, N. V. Zamanov, Yu. K. Mironov, V. S. Romanov, S. M. Maznichenko